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# DFT Study in Strain Engineering HfSe<sub>2</sub> Nanosheets for Sensing SO<sub>2</sub>, SOF<sub>2</sub>, and SO<sub>2</sub>F<sub>2</sub> Gases

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Abstract-Strain engineering is a technique of tuning a material's properties by modifying its mechanical or structural characteristics. Using first-principles calculations, the strain impact on the electronic structure of layered HfSe<sub>2</sub> nanosheets has been studied to examine the sensing capabilities of the layered HfSe<sub>2</sub> nanosheets to three SF<sub>6</sub> decomposition compounds (SO<sub>2</sub>, SO<sub>2</sub>F<sub>2</sub>, and SOF<sub>2</sub>). With biaxial strain, the band gap of the surfaces under discussion expands with tensile up to 1% and gradually decreases for strains between -1% and -5%. Additionally, our findings demonstrate that biaxial strains have a significant impact on the electronic characteristics, adsorption energy (E<sub>ad</sub>), and desorption time of SO<sub>2</sub>, SOF<sub>2</sub>, and SO<sub>2</sub>F<sub>2</sub> adsorbed HfSe<sub>2</sub> systems. The highest adsorption impact of SO<sub>2</sub>, SOF<sub>2</sub>, and SO<sub>2</sub>F<sub>2</sub> on the bilayer of HfSe<sub>2</sub> occurred at a compressive strain of -5%, with the E<sub>ad</sub> determined as -1.34, -1.27, and -1.31 eV, respectively. These external factors are widely preferred and offer a workable solution for tunable HfSe2-based gas sensors and electrical devices under biaxial compressive strain.

Keywords— Strain engineering, Monolayer HfSe<sub>2</sub>, Bilayer HfSe<sub>2</sub>, SF<sub>6</sub> gas, Gas sensor.

## I. INTRODUCTION

To ensure the electricity system's safe operation,  $SF_6$ gas, which has exceptional insulating strength and arcextinction properties, is frequently used in extra-high voltage equipment [1-2]. However, the inevitably occurring insulating flaws, such as partial discharge and partial overheating during extended use of power equipment, would break down SF<sub>6</sub> over time into several ionized F atoms and low-fluorine sulfides [3]. While these low-fluorine sulfides can often recombine with F atoms to produce SF<sub>6</sub> once more, in the presence of trace amounts of water and oxygen, the sulfides would continue to react with the impure elements to generate a variety of stable gas species, including H<sub>2</sub>S, SO<sub>2</sub>,  $SOF_2$ , and  $SO_2F_2$  [4-5]. These  $SF_6$  decomposition components are dangerous to the environment and human health, speeding up facility corrosion and raising the likelihood of system paralysis [6]. Therefore, operational detection of these typical  $SF_6$  breakdown components is crucial. In this regard, it has been determined that the

detection of  $SF_6$  decomposed species, particularly  $SO_2$ ,  $SOF_2$ , and  $SO_2F_2$ , is a practical and efficient method for assessing the operational state of  $SF_6$  insulating devices, which is also significant to ensure the smooth operation of the entire power system [7-8]. Therefore, monitoring the decomposition components, taking precautions promptly, and ensuring the safe operation of power equipment would be possible and practical using gas sensors and adsorbents with ultrahigh-sensitivity and selective qualities [4].

As a result, a wide range of two-dimensional (2D) layered nanosheets, including graphene [9], boronphosphorus [10], aluminium nitrate [11], aluminum phosphate [12], and silicone [13]. They draw considerable interest due to their distinctive architectures and exceptional characteristics. Due to the structural stability, chemical inertness, and physical properties, the 2D MX<sub>2</sub> family (transition metal dichalcogenides (TMDs); MX<sub>2</sub>: M = transition metals (TM); X = chalcogen) has lately demonstrated tremendous potential as gas sensors [14-16]. TMD materials for gas sensing have received a lot of attention, including MoS<sub>2</sub> [17, 18], MoSe<sub>2</sub> [19], PdS<sub>2</sub> [16],  $PtTe_2$  [20], and  $wS_2$  [21]. For example, Li et al. [22] examine and evaluate the outcomes of  $MoS_2$ -based gas sensors and concluding that 2D MoS<sub>2</sub> may be suitable for creating high-performance ambient-temperature gas sensors. While other TMD types, such as  $ZrX_2$  and  $HfX_2$  (X = Se, S, and Te), have received less research. HfSe<sub>2</sub> nanosheet has been predicted theoretically to have a small band gap and strong carrier mobility [23]- Numerous scientific experiments have successfully produced the monolayer [24], which suggests that it might be used as a material for gas sensors.

A popular technique for enhancing the nano-system's chemical reaction and catalytic behaviour, and hence enhancing the adsorption and sensing capabilities of the nano-materials, is surface-doping by TM [25-29]. Extensive research on TMDs as  $SF_6$  decomposition sensors, such as  $MoS_2[30, 31]$  and  $MoTe_2$  [32]. sensors, has been done during the last few decades. According to the investigations, Ga- $MoS_2$  might be an excellent gas-sensing material for SO<sub>2</sub> and

This work is licensed under a <u>Creative Commons Attribution 4.0 International License</u>. https://doi.org/10.32792/utq/utjsci/v12i1.1393 SO<sub>2</sub>F<sub>2</sub>; however, because of its limited adsorption, it was not a good fit for SOF<sub>2</sub> sensing [31]. On the other hand, using first-principles calculations, Cui Hao et al. [33-34] investigated the adsorption behaviour of pristine, and Ptdoped HfSe<sub>2</sub> monolayers to some hazardous gases (SOF<sub>2</sub>,  $SO_2$ , and  $NO_2$ ). It is discovered that the Pt-doped HfSe<sub>2</sub> monolayer exhibits better behavior for SO<sub>2</sub>. Rh-doped HfSe<sub>2</sub>'s adsorption nature and behavior were investigated by Wang et al. [35], and their findings indicate that SO<sub>2</sub> and the monolayer interact more strongly than SO<sub>2</sub>F<sub>2</sub>. Additionally, our group [36] reported that monolayer HfSe<sub>2</sub> doped with a nickel cluster could distinguish SF<sub>6</sub> decomposition gases, also exhibiting excellent sensitivity and regulated gas release/ storage. Besides, nano-sensors improved sensing capabilities can enable their use in harsher environments with a strong sensitivity toward gaseous species [37-38], which is urgently needed in material science and engineering [39-40]. For adsorption systems, control over charge transfer and electronic structures would impact both sensing performance and adsorption capabilities [41-42].

Applying external strain to the adsorption system is one efficient method for tuning the electronic characteristics [43- 44]. For example, Guzman and Strachen [45] forecast significant changes in bandgap for several monolayer TMDs with biaxial strain ranging from -5% to 5%. According to Li et al.'s research [46] The Ru-PtSe<sub>2</sub> monolayer has a tunable sensing response for gas detection with modulated strains, as evidenced by its increased sensitivity to C<sub>2</sub>H<sub>2</sub> detection by manipulating biaxial strains. The SnS<sub>2</sub> sheets also demonstrate how biaxial strains can significantly alter the adsorption energy, electronic characteristics, and charge transfer of systems with adsorbed nitrogen dioxide [47]. According to Ni et al.'s findings, the adsorption capabilities of an Ag-doped WSe<sub>2</sub> monolayer declined as the biaxial strain increased [48].

In this study, we design models for  $SO_2$ ,  $SOF_2$ , and  $SO_2F_2$  adsorption on pure monolayer and bilayer  $HfSe_2$  nanosheets to compute the adsorption behaviors, and investigate their potential as gas sensors under applying biaxial strain. Our research may pave the path for using  $HfSe_2$  to identify  $SF_6$  decompositions in challenging application environments, including the interior of gasinsulated switchgear (GIS) equipment.

## II. COMPUTATIONAL METHODS

All computations on HfSe<sub>2</sub> nanosheets are performed with DFT via the DMol3 package[49] by using the material studio program, and considering the impact of spin polarization, peripheral electrons for Hf (Se) are produced in a  $5d^2 6s^2 (4s^24p^4)$  configuration (32 Se atoms and 16 Hf atoms). The generalized approximation gradient (GGA) and Perdew-Burke-Ernzerhof (PBE) are utilized to handle the exchange-related energy [50], where the PBE functional is very popular because it is a non-empirical functional with reasonable accuracy over a wide range of systems. While PBE is typically not the most accurate GGA functional (especially for band gap) for a given system, it usually is not too far off either [51]. In addition, GGA struggles to adequately describe long-range effects like vdWs [52]. To investigate the reactions between HfSe<sub>2</sub> sheets and gas molecules, the long-range dispersion-corrected developed by

Tkatchenko-Scheffler (TS) has been utilized. TS offers empirical parameters for the elements' surface (HfSe<sub>2</sub> sheets), whereas Grimme don't offered it [53]. The basis set was specificed by double numerical plus polarization (DNP). We used the appropriate number of Monkhorst-Pack k-point grids,  $7 \times 7 \times 1$ , and mesh smearing of 0.005 Ha for the geometry optimization and the electronic attribute computations. All structural models were completely relaxed until the maximum force and the tolerance energy were less than 0.002 Ha/Å and 10<sup>-5</sup> eV, respectively. In addition, the self-consistent field (SCF) tolerance was set to 10<sup>-6</sup> Ha [54]. A supercell was large enough for  $4 \times 4 \times 1$  HfSe<sub>2</sub> sheets [33, 35, 55]. To avoid erroneous interactions between the periodic pictures, a 20 vacuum was placed between them along the zaxis. The adsorption energy  $(E_{ad})$  is defined as the energy required to adsorb a certain amount of gas onto a HfSe<sub>2</sub> sheet [56], which is quantitatively described as follows:

$$E_{ad} = E_{surf/gas} - E_{surf} - E_{gas} \tag{1}$$

where  $E_{surf/gas}$ ,  $E_{surf}$  and  $E_{gas}$  are the total energy of gas adsorption systems, isolated surfaces, and the separated gas molecules, respectively. Here, the stronger the adsorption, the more negative the  $E_{ads}$ . Hirshfeld population analysis can be used to examine the charge transfer (Q<sub>t</sub>) of the gas molecules adsorbed on layered HfSe<sub>2</sub> sheets [56]. Q<sub>t</sub> denotes the quantity of charge transfer between substrates and the gas molecules.

$$Q_t = Q_{tiso} - Q_{tads} \tag{2}$$

where  $Q_{tiso}$  stands for the isolated gas molecules and  $Q_{tads}$  for the total charges carried of the adsorbed gas.

## III. RESULTS AND DISCUSSION

## A. The Geometric and Electronic Properties

Fig. 1 (a, b) shows the top and side views of the optimized structures of the HfSe<sub>2</sub> monolayer and bilayer. The P3-m1 space group (No. 164) is home to the tetragonal symmetry found in HfSe<sub>2</sub> sheets. Se, Hf, and Se atoms are stacked one on top of the other to form their structure. The internal coordinates are (0, 0, 0) and (1/3, 2/3, 0.227) for Hf and Se, respectively. The researched material's optimal lattice parameters are a = b = 3.75 Å, which agree with the theoretical conclusions that have already been established [34, 36]. Interatomic distance Se-Hf is 2.69 Å at equilibrium, and the lattice angles of  $\alpha = \beta = 90.00^{\circ}$  and  $\gamma = 120.00^{\circ}$ . In our study, AA (eclipsed with Se (Hf) over Se (Hf)) is the order in which the HfSe<sub>2</sub> bilayer was stacked; the interlayer gap between Se and Se is 6.084 Å. The band structure (BS) and density of states (DOS) of the unstrained HfSe<sub>2</sub> monolaver and bilayer are shown in Fig. 1(c, d), respectively. Given that the conduction band (CB) minimum is identified at the M point and the valence band (VB) maximum is discovered at point  $\Gamma$  close to the Fermi level, the profile of the derived BS suggests that the investigated sheets are an indirect semiconductor. The monolayer and bilayer systems' resulting band gaps are 0.477 eV (which agrees with the previously reported value [34-54]) and 0.429 eV, respectively. As observed in the figure, the extent of the energies is limited from -3 eV to 3 eV (-2 eV to 2 eV) for monolayer (bilayer), the VB is dominated mainly by Se atoms, with a modest contribution from Hf atoms, while both atoms contribute to the CB. The response of Se and Hf atoms in the HfSe<sub>2</sub> monolayer and bilayer is reasonable for the hybridization among Se 4p and Hf 5d states, and the hybridization occurs

significantly farther from the Fermi level.



Fig. 1: Geometric structure, BS, and DOS of (a,c) monolayer HfSe<sub>2</sub> and (b,d) bilayer HfSe<sub>2</sub>.

## B. Adsorption Unstrained Systems

To determine the most favorable adsorption configuration, the three gas molecules were placed molecules at three possible sites above the pristine  $HfSe_2$  monolayer and bilayer positions classified as: 1-the top Se atom of the downer layer, i.e. gas molecules above the center of hexagonal ring of  $HfSe_2$ , 2- the top of Hf atom, 3- the top of Se of the upper layer. The three aim gases were adsorbed using the best model of  $HfSe_2$  sheets. The adsorption features of  $HfSe_2$  sheets were investigated in light of the adsorption parameters (represented by  $E_{ad}$ ,  $Q_t$ , and distance between Se and S atoms), which are listed in Table 1. The ideal adsorption models of three gases (SO<sub>2</sub>, SOF<sub>2</sub>, and SO<sub>2</sub>F<sub>2</sub>) on the  $HfSe_2$  sheets are shown in Fig. 2. All systems have long adsorption lengths and low  $Q_t$  value, indicating insensitive physisorption, as indicated by the values in Table 1.

 $\label{eq:source} \begin{array}{c} \textbf{Table 1.} & A dsorption \ characteristics \ of \ SO_2, \ SOF_2, \ and \ SO_2F_2 \ gases \ on \\ & \ HfSe_2 \ sheets \end{array}$ 

Structure	Gas molecules	Figure	E <sub>ad</sub> (eV)	Q <sub>t</sub> Hirshfel d (e)	D (Å) S-Se
HfSe <sub>2</sub> monolayer	$SO_2$	Fig. 2(a1)	-0.32	-0.076	3.194
	SOF <sub>2</sub>	Fig. 2(b1)	-0.247	-0.0008	3.692
	$SO_2F_2$	Fig. 2(c1)	-0.279	-0.0031	3.866
HfSe <sub>2</sub> bilayer	$SO_2$	Fig. 2(a2)	-0.294	-0.0785	3.173
	SOF <sub>2</sub>	Fig. 2(b2)	-0.214	-0.0015	3.687
	$SO_2F_2$	Fig. 2(c2)	-0.255	-0.0047	4.03



**Fig.2:** Side view of SO<sub>2</sub>, SOF<sub>2</sub> and SO<sub>2</sub>F<sub>2</sub> adsorption systems, above (a1-a3) monolayer and (b1-b2) bilayer HfSe<sub>2</sub>.

#### C. Electronic Properties under Strain

Fig. 1 shows the schematic diagrams of applying tensile and compressive biaxial strain on  $HfSe_2$  sheets before gas adsorption, where positive and negative values represent tensile and compressive strain, respectively. Strain engineering is one of the most efficient ways to control the electronic characteristics of semiconductors. It is widely known that altering the electronic characteristics of semiconductor materials is crucial for gas sensors. Biaxial strain is more effective than uniaxial strain in adjusting the adsorption properties [57]. Based on the literature that predicted large band gap changes for several monolayer TMDs with biaxial strain ranging from -5% to 1%. The monolayer shows a transition from semiconductor to metal at -5% with surface stability [45]. Therefore, we examined the electronic characteristics of the monolayer and bilayer of HfSe<sub>2</sub> while subjected to an external biaxial strain. The lattice constants of the strained and unstrained systems, a and  $a_0$ , respectively, are used to calculate the biaxial strain [56]. which is defined as  $\varepsilon = (a - a_0)/a_0 \times 100\%$ . Only the atomic position can relax for the strained system. Here, strains between -5% and 1% with a step of 2% are taken into consideration; the strain is set with lattice constants of 3.563, and 3.713, 3.788 respectively. Three 3.637, Å, negative/positive strain values correspond to compressive/tensile strains. Analyzing their BSs and DOS prior to gas adsorption is vital to comprehend their adsorption properties under biaxial strain with values -5%, -3%, -1%, and 1%, see Figs. 3 and 4 of monolayer and bilayer, respectively. We discovered that the monolayer and the bilayer have an indirect BG at zero strain and that the gap (decreases) increases at (compressive) tensile uniform strain. The band gaps of the monolayer and bilayer remain indirect for all strain values considered. In the case of a bilayer, the maximum VB changes from  $\Gamma$  to M because the state at M gains energy under tensile strain, although the corresponding eigenvalue at  $\Gamma$  does not significantly change. We can observe that the minimum CB and the maximum VB were always found at the M and  $\Gamma$  point, respectively. Hence, the lack of contrast in maximum VB and the minimum CB governed the non-switch from indirect to direct BG. Keep in mind that negative strain values represent compressive strain. The impact of mechanical distortions on the electronic structure is primarily discernible at the M-K path, where the bottom of the CB is drastically decreased relative to the equilibrium status. The top of VB at the  $\Gamma$  spot travels upward concurrently. Because of the alteration in bond lengths, this behavior is understandable makes sense because the orbital overlap between the Hf and Se atoms reduces. The electron conductivity in semiconducting TMD nanosheets can be increased or decreased due to changes in the electronic architecture caused by mechanical deformations. The BG is decreased due to the 2D-compressive strain, and the electron conductivity at the Fermi level is noticeably increased. Since the BGs of the HfSe<sub>2</sub> monolayer (bilayer) vanish at about -5% (-3%) strain, the transport paths are open, enabling electrical conductivity. This can be described by the Brillouin zone deformation, where the applied strain evenly alters the primitive cell, bands penetrate the Fermi level, and as a result, a monolayer's (bilayer's) character transforms from semiconducting to metallic. Additionally, it is intriguing to observe that the bilayer's interlayer spacing marginally increases with compressive strain and reduces with tensile strain, as seen in Fig. 5.

#### D. Enhancement of sensing features under strain

Adsorption systems under biaxial strain ranging from  $-5\% \sim$ 1% are monitored in order to assess the impact of applied strain on the potential for adsorption and sensing performance of HfSe<sub>2</sub> sheets towards SO<sub>2</sub>, SOF<sub>2</sub>, and SO<sub>2</sub>F<sub>2</sub>. The E<sub>ads</sub> and Q<sub>t</sub> with various biaxial strains are depicted in Fig. 6 and 7. The findings demonstrate that applying biaxial strain significantly affects the Ead for all systems, increasing E<sub>ads</sub> as tensile strain increases. However, as the compressive strain is increased, the adsorption diminishes. In both systems, a better response performance for the sensing material to the target gas is generally correlated with lower  $E_{ads}.$  The  $SO_2F_2$  molecule has the same behavior in both systems, but the bilayer exhibits a more considerable sensitivity to the  $SO_2$  and  $SOF_2$  molecules under compressive strains of up to -5% as the  $E_{ad}$  gets more negative (-1.34 and -1.27 eV, respectively).



Fig. 3: BGs and DOS for monolayer HfSe2 in the presence of biaxial strain, represented by the dashed red line at the Fermi level (tuning to zero).



Fig. 4: BGs and DOS for bilayer HfSe2 in the presence of biaxial strain represented by the dashed red line at the Fermi level (tuning to zero).



Fig. 5: Interlayer distance for bilayer  $HfSe_2$  as a function of biaxial strain applied, distance in Å.

On the other hand, when the tensile strain reaches 1%, it is expected that the  $E_{ad}$  of the HfSe<sub>2</sub> bilayer system will be -0.0803, + 0.0038, and -0.0403 eV for SO<sub>2</sub>, SOF<sub>2</sub>, and SO<sub>2</sub>F<sub>2</sub>, respectively. This indicates that the adsorption mechanism is weak or non-existent for bilayer systems with a 1% tensile strain. Due to the bilayer system's instability at 1% tensile biaxial strain, the gas molecules are independent of the surface. In other words, the system had some difficulty in adsorption of these gases[31].



Fig. 6: Variation of the  $E_{ads}$  for SO<sub>2</sub>, SOF<sub>2</sub>, and SO<sub>2</sub>F<sub>2</sub> adsorbed on (a) monolayer HfSe<sub>2</sub> and (b) bilayer HfSe<sub>2</sub> as a function of biaxial strain.

In related adsorption systems, the  $Q_t$  exhibits a similar tendency as a function of biaxial strain. After adjusting the biaxial strain from -5% to 1%, -0.080 and -0.0787 electrons were transferred by the HfSe<sub>2</sub> monolayer and bilayer from the SO<sub>2</sub> molecule at -5% strain, respectively. Upon examining different systems, it is notable that SO<sub>2</sub>-HfSe<sub>2</sub> adsorption systems have the highest  $Q_t$  sensitivity. After applying the biaxial strain, there was clear change in the  $Q_t$  of the SO<sub>2</sub>-HfSe<sub>2</sub> system from -0.08 to -0.0903e for biaxial strain from -5% to 1%, respectively. This result is considered important compared with the charge transfer values in Ref. [58]. where  $Q_t$  ranges from 0.008 to 0.019e for biaxial strain from -4% to 1%, respectively.



Fig. 7: Variation of the  $Q_t$  for SO<sub>2</sub>, SOF<sub>2</sub>, and SO<sub>2</sub>F<sub>2</sub> adsorbed on (a) monolayer HfSe<sub>2</sub> and (b) bilayer HfSe<sub>2</sub> as a function of biaxial strain.

In contrast, the SOF<sub>2</sub> and SO<sub>2</sub>F<sub>2</sub> adsorption systems (related to low charge transfer) do not experience any substantial Q<sub>t</sub> following the addition of external strain. According to Hirshfeld analysis, for SOF<sub>2</sub>-adsorbed monolayer and bilayer HfSe<sub>2</sub> nanosheet at -5% strain, there are 0.0027 e and 0.003 e moving from the substrate to SOF<sub>2</sub> molecule, respectively. The SO<sub>2</sub>F<sub>2</sub> adsorption systems behave similarly to SOF<sub>2</sub> systems in Q<sub>t</sub> as well. Even if these two gases continue to have the electron donor feature, we might see a change in the charge transport path in the two systems when the applied strain drops below -3%.

In Fig. 8(a and b), we present the relationship between strain and the adsorption distance (d) of HfSe<sub>2</sub> systems. Both systems show slight decrease in the S-Se adsorption distance for the gas molecules with the increase of compression strain (from 0 to -5%), where the absolute value of  $E_{ad}$  increases from 0.32 ~ 0.86 eV (0.29 ~ 1.34 eV), 0.247~0.85 (0.22 ~ 1.27 eV)and 0.279~1.3 (0.255 ~ 1.31 eV) for the adsorption SO<sub>2</sub>, SOF<sub>2</sub>, and SO<sub>2</sub>F<sub>2</sub> on HfSe<sub>2</sub> monolayer (bilayer), and a decreasing tendency in the d-curve indicates stronger binding at a reduced adsorption distance. Except for the adsorption of SO<sub>2</sub>F<sub>2</sub> gas on the bilayer, we can see that the distance increases and then decreases at -5% strain.



**Fig. 8:** Variation of the adsorption distance for SO<sub>2</sub>, SOF<sub>2</sub>, and SO<sub>2</sub>F<sub>2</sub> adsorbed on (a)  $HfSe_2$  monolayer and (b)  $HfSe_2$  bilayer as a function of biaxial strain.

Initially, we computed the BSs of the studied structures under biaxial strain. The variance in BG with strain is illustrated in Fig. 3 and 4. Generally, compressive strain causes a decrease in the BG of all studied systems. The effect of tensile was slightly different, the BG rising as the biaxial strain increased. The BGs of the gas-adsorbed systems exhibit a linear relationship in Fig. 9 (a and b), slightly narrower when subjected to compressive strain and broader when tensile strains are applied. In other words, despite both systems being strain-sensitive and giving diverse sensing responses, the HfSe<sub>2</sub> bilayer has greater sensitivity upon gas detection with the modification of biaxial stresses. Adding biaxial strain has less impact the unadsorbed HfSe<sub>2</sub> monolayer, which retains its semiconductor nature except for a compressive strain of up to -5% when the monolayer turns metallic.



**Fig. 9:** Variation of the BS for SO<sub>2</sub>, SOF<sub>2</sub>, and SO<sub>2</sub>F<sub>2</sub> adsorbed on (a)  $HfSe_2$  monolayer and (b)  $HfSe_2$  bilayer as a function of biaxial strain.

#### E. Recovery Time

Recovery time ( $\tau$ ), which indicates whether the material is acceptable for practice application reality or not, is another essential merit quality. According to the Arrhenius formula, the relationship between  $\tau$  and desorption barrier  $E_a$ (is equal to  $E_{ad}$ ) is estimated as follows:

$$= v_0^{-1} e^{-E_a/k_B T}$$
(3)

Here,  $v_0$  is the operational frequency, which is taken as  $10^{12}$ Hz,  $k_B$  is the Boltzmann constant, and T is the operational temperature (300 K) [59]. The  $\tau$  is defined by the time it takes the target material (HfSe<sub>2</sub> sheets) to reestablish its pristine properties or condition once the gas molecule is removed, determining how long the HfSe<sub>2</sub> sheets can be reused. Fig. 10(a, b) shows the weighted  $\tau$  as a function of biaxial strain. According to the calculations shown in Fig. 10, the adsorption time of gas molecules of all systems gets faster (or instantaneously) under the tensile strains and increases linearly under the compressive strains.

For the adsorption on the HfSe<sub>2</sub> monolayer, SO<sub>2</sub>F<sub>2</sub> has the largest  $E_{ad}$  under the compressive strain of up -5%, which means that the reaction between SO<sub>2</sub>F<sub>2</sub> and the HfSe<sub>2</sub> monolayer is the most powerful. However, the desorption barrier is higher with a high  $E_{ad}$ . Since the SO<sub>2</sub>F<sub>2</sub> molecule has a much higher estimate of the  $\tau$  (9.8×10<sup>9</sup> s) than the other two gas molecules, it has the lowest absorption rate. The other two gases have a higher absorption rate due to a lower absorption barrier, and the corresponding  $\tau$  (6.6 and 5.15 min for SO<sub>2</sub> and SOF<sub>2</sub> gases, respectively) is also shorter.

For the desorption from the HfSe<sub>2</sub> bilayer, all the  $E_{ads}$  of these three gas molecules are more significant than those on the HfSe<sub>2</sub> monolayer under the compressive strain. They have higher desorption barriers as a result, which results in slower absorption rates and longer  $\tau$ . Applying compressive strain significantly lengthens the  $\tau$ . HfSe<sub>2</sub> sheets have a

more remarkable detection ability when the target gas binds to the surface for a long time. On the other hand, biaxial tensile strain would result in a noticeably reduced  $\tau$  for each system under consideration, indicating that biaxial tensile strain might be used to desorb gas quickly. For instance, the  $\tau$  of the gases/HfSe<sub>2</sub> bilayer adsorption system is drastically reduced when the external strain of -5% changes to that of 1%. As a result, the analysis of  $\tau$  shows that controlling external factors is necessary to regulate the adsorption-absorption process. Due to the combination of the sheets' selectivity of HfSe<sub>2</sub>, greater  $E_{ad}$ , and the high  $\tau$ , the bilayer is more effective than the monolayer in detecting SF<sub>6</sub> breakdown products.



Fig. 10: Variation of the recovery time for  $SO_2$ ,  $SOF_2$ , and  $SO_2F_2$  adsorbed on (a)  $HfSe_2$  monolayer and (b)  $HfSe_2$  bilayer as a function of biaxial strain.

#### IV. CONCLUSIONS

In brief, the sensor mechanism of layered HfSe<sub>2</sub> towards three common gases resulting from SF<sub>6</sub> decomposition is investigated using first-principles calculations. The geometrical stability and electronic characteristics of SO<sub>2</sub>, SOF<sub>2</sub>, and SO<sub>2</sub>F<sub>2</sub> on the HfSe<sub>2</sub> monolayer and bilayer are monitored under biaxial strain engineering. According to the electronic systems study, biaxial strain could be used to successfully control the electrical conductivities of the surfaces under consideration, while the adsorption of gas molecules does not generate noticeable changes in the electronic properties of the surfaces. However, the Ead of all systems increases after applying compressive strain, while the tensile biaxial strain decreases the  $E_{ad}.$  The findings demonstrate that the  $E_{ads}$  of SO<sub>2</sub> and SOF<sub>2</sub> molecules on a monolayer (-0.86 and -0.85 eV, respectively) are lower than that for SO<sub>2</sub>F<sub>2</sub> adsorption (-1.30 eV) system at -5% strain, indicating that the application of compressive strain increased the HfSe<sub>2</sub> monolayer's sensitivity to  $SO_2F_2$ . In addition, the bilayer shows high adsorption energy towards all gases in the following order:  $SO_2(-1.34) > SO_2F_2(-1.31) > SOF_2(-1.27)$ . Moreover, the computed recovery times for each system showed that tensile strain assisted-absorption, whereas compressive strain enhanced gas adsorption. The study suggested that HfSe<sub>2</sub> nanosheets are a promising bidirectional sensing material for SF<sub>6</sub> decomposition detection, which can be manipulated by applying biaxial strain.

#### CONFLICT OF INTEREST

Authors declare that they have no conflict of interest.

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